

Canonical quantization:

Take 3-manifold M_3 to be of the form

$$M_3 \cong \Sigma \times \mathbb{R}$$

↑
Riemann surface

Canonical quantization \rightarrow Hilbert space \mathcal{H}_Σ

Recall:

$$CS(A) = \frac{\kappa}{8\pi^2} \int_{M_3} \text{Tr} \left(A \wedge dA + \frac{2}{3} A \wedge A \wedge A \right)$$

Taking $M_3 = \Sigma \times \mathbb{R}$ and choosing the gauge $A_0 = 0$, gives

$$\mathcal{L} = \frac{\kappa}{8\pi^2} \int dt \int_{\Sigma} \varepsilon^{ij} \text{Tr} A_i \frac{d}{dt} A_j$$

\rightarrow Poisson brackets :

$$\{A_i^a(x), A_j^b(y)\} = \frac{4\pi^2}{\kappa} \cdot \varepsilon_{ij} \delta^{ab} \delta^2(x-y)$$

Due to gauge choice, system is subject to:

$$\frac{\delta \mathcal{L}}{\delta A_0} = 0 \iff \varepsilon^{ij} F_{ij}^a = 0 \quad (*)$$

We quantize the system by first constraining the phase space \rightarrow symplectic quotient:

recall $\mathcal{A}_\Sigma \cong \Omega'(\Sigma, \mathfrak{g})$ symplectic manifold with non-degenerate symp. form

$$\omega(\alpha, \beta) = -\frac{\kappa}{8\pi^2} \int_\Sigma \text{Tr}(\alpha \wedge \beta), \quad \alpha, \beta \in \Omega'(\Sigma, \mathfrak{g})$$

Then $\mathcal{G}_\Sigma \cong \text{Map}(\Sigma, G)$ acts on \mathcal{A}_Σ

by gauge transformations

Imposing the constraint (*) is then equivalent to defining the quotient space:

$$\mathcal{M}_\Sigma = F^{-1}(0) / \mathcal{G}_\Sigma$$

"symplectic quotient"

The symplectic structure on \mathcal{M}_Σ is given by

$$\omega_A(\tilde{\alpha}, \tilde{\beta}) = \omega_A(\alpha, \beta) \quad \text{for } \tilde{\alpha}, \tilde{\beta} \in T_A \mathcal{M}_\Sigma$$

\rightarrow definition is independent of choice of

$$A, \alpha, \beta$$

Holomorphic viewpoint:

Lift action of gauge group G_Σ to L_Σ .

→ complex line bundle \mathcal{L} on M_Σ
with connection ∇ of curvature $-i\omega$

To quantize M_Σ , we need to pick hol. sections of \mathcal{L} → choice of complex structure

1) pick a complex structure \mathcal{J} on Σ

2) for $\alpha \in T^*\Sigma$, define

$$\mathcal{J}_{T^*\Sigma} \alpha = -\mathcal{J} \alpha$$

3) $T^*\Sigma = T^{(1,0)}\Sigma \oplus T^{(0,1)}\Sigma$,

where $T^{(1,0)}\Sigma$ and $T^{(0,1)}\Sigma$ consist of $(0,1)$ -forms and $(1,0)$ -forms on Σ with values in $\mathfrak{ad}(\mathbb{Q})$.

4) ω is of type $(1,1)$.

5) the connection d_A decomposes as

$$d_A = \partial_A + \bar{\partial}_A,$$

6) complexified $G_{\mathbb{C}}$ -action is determined by:

$$\bar{\partial}_A \mapsto g \cdot \bar{\partial}_A \cdot g^{-1}, \quad \partial_A \mapsto \bar{g} \cdot \partial_A \cdot \bar{g}^{-1}$$

7) The set $\mathcal{A}/G_{\mathbb{C}}$ can be identified with the set $\mathcal{M}_{\mathbb{C}}$ of equivalence classes of hol. structures on $\mathbb{C}P^1$.

8) $\mathcal{T} \equiv$ moduli space of complex structures on Σ

$\rightarrow t \in \mathcal{T}$ determines \mathcal{J}_t and $\mathcal{M}_{\mathcal{J}_t}$

$\rightarrow \mathcal{M} \times \mathcal{T}$ can be regarded as a bundle over \mathcal{T} with fibers varying with t .

9) Introduce quantum bundle $\tilde{\mathcal{H}}_{\mathbb{C}}$ over \mathcal{T} . The fiber $\tilde{\mathcal{H}}_{\mathbb{C}}|_t$ over a point t is $H^0(\mathcal{M}_{\mathcal{J}_t}, \mathcal{L})$ (holomorphic sections)

10) There is a "projectively flat conn." ∇ on $\tilde{\mathcal{H}}_{\mathbb{C}} \rightarrow \mathcal{T}$.

Construction of the connection

First, we need to describe the connection ∇ on L_Σ , with curvature $-i\omega$.

Define
$$\nabla_u \psi(A) = \int_\Sigma d^2z u_z^a \frac{D}{DA_z^a(z)} \psi(A),$$

$$\nabla_{\bar{u}} \psi(A) = \int_\Sigma d^2z \bar{u}_{\bar{z}}^a \frac{D}{DA_{\bar{z}}^a(z)} \psi(A),$$

for u and \bar{u} adjoint valued $(1,0)$ and $(0,1)$ forms on Σ and $d^2z = idz d\bar{z}$.

We need

$$\left[\frac{D}{DA_\omega^a(\omega)}, \frac{D}{DA_{\bar{z}}^b(z)} \right] = \overbrace{-i \frac{k}{4\pi} \delta_{ab} \delta_{z\bar{z}}(z, \omega)} = 2\pi i \omega$$

$$\left[\frac{D}{DA_z^a(z)}, \frac{D}{DA_\omega^b(\omega)} \right] = \left[\frac{D}{DA_{\bar{z}}^a(z)}, \frac{D}{DA_{\bar{\omega}}^b(\omega)} \right] = 0$$

→ take:

$$\frac{D}{DA_z} = \frac{\delta}{\delta A_z} - \frac{k}{8\pi} A_{\bar{z}}$$

$$\frac{D}{DA_{\bar{z}}} = \frac{\delta}{\delta A_{\bar{z}}} + \frac{k}{8\pi} A_z$$

→ holomorphic sections have to satisfy:

$$\frac{D}{DA_z} \Psi = 0, \quad \Psi \in H^0(\mathcal{A}_Z, L_Z) =: \mathcal{H}_G|_Y$$

To study the objects $\tilde{\mathcal{H}}_G|_Y = H^0(\mathcal{M}_Z, \mathcal{L})$, we need to further restrict to $\mathfrak{g}_\mathbb{C}$ -inv. subspace of $\mathcal{H}_G|_Y$.

→ the condition for $\mathfrak{g}_\mathbb{C}$ -invariance is

$$\left(-D_{\bar{z}} \frac{D}{DA_{\bar{z}}^a(z)} + \frac{\kappa}{4\pi} \delta_{ab} F_{\bar{z}z}^b(z) \right) \Psi = 0$$

where $F_{\bar{z}z}^a(z)$ is the curvature of the connection A , and $D_{\bar{z}} = \frac{\partial}{\partial \bar{z}} + A_{\bar{z}}$.

We will now construct the connection on

$$\begin{array}{ccc} \tilde{\mathcal{H}}_G|_Y & \longrightarrow & \mathcal{V} \\ & & \downarrow \\ & & \mathcal{F} \end{array}$$

Note that the space \mathcal{A}_Z is infinite-dim., while \mathcal{M}_Z has complex dimension

$$\dim_{\mathbb{C}} \mathcal{M}_Z = (g-1) \cdot \dim G$$

This comes from the fact that flat connections are characterized by "Wilson lines" (holonomies around cycles of Σ)

\mathcal{T} is a complex manifold
 \rightarrow write $g^{(1,0)}$ and $g^{(0,1)}$ respectively for the ∂ and $\bar{\partial}$ operators on \mathcal{T} .

Explicitly :

$$g^{(1,0)} = \int_{\Sigma} \delta t_{\bar{z}\bar{z}} \frac{\delta}{\delta t_{\bar{z}\bar{z}}}$$

$$g^{(0,1)} = \int_{\Sigma} \delta t_{zz} \frac{\delta}{\delta t_{zz}}$$

Working "upstairs" on \mathcal{A} , the projectively flat connection on $\mathcal{H}\mathbb{Q}$ is

$$g^{\mathcal{H}\mathbb{Q}} = g^{(1,0)} - \frac{it}{4} \cdot \frac{4\pi}{k} \int_{\Sigma} \delta \gamma^{\bar{z}} \frac{D}{DA_z^a(z)} \frac{D}{DA_{\bar{z}}^a(z)}$$